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A multiwavelet Galerkin boundary element method for the stationary Stokes problem in 3D

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ABSTRACT

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1. Introduction

The Stokes problem has been usually applied to model incompressible creeping flows where the fluid Reynolds number is very low. The problem has been studied extensively by boundary element method (BEM). The BEM is often more convenient than the traditional finite element method (FEM), since the corresponding equations are formulated on the boundary, which satisfy the incompressibility constraint. However, the system matrices of boundary element equations are computed densely due to the nonlocal nature of the boundary integral operators, which is the main disadvantage of the BEM compared with FEM that leads to sparse matrices. This drawback makes inconvenience to apply the BEM to the large-scale problems. However, many methods for the fast solution of BEM have been developed in the last decades. Prominent examples for such methods are the fast multipole method [1–3], the multi-level BEM [4,5], the panel clustering [6], H-matrices [7] and the adaptive cross approximation [8]. Furthermore, wavelet Galerkin BEM (GBEM) [9-29] was introduced and successfully applied to many practical problems in the last years.

In Ref. [29], we have described a GBEM using Alpert multiwavelets proposed in Ref. [10] for solving the two-dimensional Stokes problem. The multiwavelets not only have short supports and high order of vanishing moments but also are simply piecewise polynomials which allow easy and efficient evaluation of the matrix entries. In this paper, we will present a multiwavelet GBEM (MGBEM) for solving the three-dimensional Stokes

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In this paper, a multiwavelet Galerkin boundary element method is presented for the fast solution of the stationary Stokes problem in three dimensions. Piecewise linear discontinuous multiwavelet bases are constructed on each patch of piecewise smooth surface individually, which allow easy and efficient evaluation of the matrix entries. Because of the use of the multiwavelets, the system matrix can be compressed to O(N) (N denotes the number of unknowns) nonzero entries without compromising the order of convergence as for the conventional Galerkin boundary element method. Numerical results of two test samples are given to demonstrate the availability of the present method.

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problem, in which the multiwavelets are constructed on each patch individually. The multiwavelets have higher order of vanishing moment compared with Haar wavelets [17]; hence, we can obtain a sparser system matrix containing O(N) nonzero entries by using two step compression strategies. What's more, the multiwavelets are simpler piecewise polynomials compared with Spline wavelets [12,18,19,22], therefore, we can calculate the matrix entries in shorter times.

We consider the stationary Stokes problem with Dirichlet boundary condition

$$\begin{cases} -\mu \Delta u_i(\mathbf{x}) + \frac{\partial p(\mathbf{x})}{\partial x_i} = 0, & \mathbf{x} \in \Omega \text{ or } \Omega', \\ \nabla \cdot \mathbf{u}(\mathbf{x}) = \sum_{i=1}^{3} \frac{\partial u_i(\mathbf{x})}{\partial x_i} = 0, & \mathbf{x} \in \Omega \text{ or } \Omega', \\ u_i(\mathbf{x}) = g_i(\mathbf{x}), & \mathbf{x} \in \Gamma, \quad i = 1, 2, 3 \end{cases}$$
(1)

where Ω is an open bounded domain in R^3 of points $\mathbf{x} = (x_1, x_2, x_3)$, its boundary Γ is assumed to be piecewise smooth, the complement of $\overline{\Omega} = \Omega + \Gamma$ is denoted by Ω' , and $\mathbf{n} = (n_1, n_2, n_3)$ denotes the unit exterior normal to Γ . The unknowns are the velocity $\mathbf{u} = (u_1, u_2, u_3)$ and pressure p of the viscous incompressible fluid filled in Ω or Ω' , here the kinematic viscosity μ is constant, $\mathbf{g} = (g_1, g_2, g_3)$ is a given function on Γ .

The solution (\mathbf{u},p) of problem (1) can be expressed in the form of the simple layer potentials [30,31,32]

$$\begin{cases} u_{j}(\mathbf{y}) = \sum_{i=1}^{3} \int_{\Gamma} t_{i}(\mathbf{x}) U_{ij}(\mathbf{x}, \mathbf{y}) dS_{\mathbf{x}}, & j = 1, 2, 3, \mathbf{y} \in \mathbf{R}^{3}, \\ p(\mathbf{y}) = \sum_{i=1}^{3} \int_{\Gamma} t_{i}(\mathbf{x}) P_{i}(\mathbf{x}, \mathbf{y}) dS_{\mathbf{x}}, & \mathbf{y} \in \mathbf{R}^{3} - \Gamma, \end{cases}$$

$$(2)$$

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