

Available online at www.sciencedirect.com

SciVerse ScienceDirect



Mathematics and Computers in Simulation 82 (2012) 984-1007

www.elsevier.com/locate/matcom

Numerical solution of the 'classical' Boussinesq system

Original article

D.C. Antonopoulos^a, V.A. Dougalis^{a,b,*}

^a Department of Mathematics, University of Athens, 15784 Zographou, Greece ^b Institute of Applied and Computational Mathematics, FORTH, 71110 Heraklion, Greece

Received 24 November 2009; received in revised form 13 September 2011; accepted 16 September 2011 Available online 18 November 2011

Abstract

We consider the 'classical' Boussinesq system of water wave theory, which belongs to the class of Boussinesq systems modelling two-way propagation of long waves of small amplitude on the surface of water in a horizontal channel. (We also consider its completely symmetric analog.) We discretize the initial-boundary-value problem for these systems, corresponding to homogeneous Dirichlet boundary conditions on the velocity variable at the endpoints of a finite interval, using fully discrete Galerkin-finite element methods of high accuracy. We use the numerical schemes as exploratory tools to study the propagation and interactions of solitary-wave solutions of these systems, as well as other properties of their solutions. © 2011 IMACS. Published by Elsevier B.V. All rights reserved.

Keywords: Water waves; 'Classical' Boussinesq systems; Initial-boundary-value problems; Fully discrete Galerkin-finite element methods; Solitary waves

1. Introduction

We consider the so-called 'classical' Boussinesq system (CB)

$$\eta_t + u_x + (\eta u)_x = 0,$$

$$u_t + \eta_x + uu_x - \frac{1}{3}u_{xxt} = 0,$$
(1.1)

for $x \in \mathbb{R}$, t > 0, supplemented by the initial conditions

$$\eta(x,0) = \eta_0(x), \qquad u(x,0) = u_0(x), \tag{1.2}$$

where η_0 , u_0 are given real functions on \mathbb{R} . The system (1.1) is a Boussinesq-type approximation of the two-dimensional Euler equations that models two-way propagation of long waves of small amplitude on the surface of an incompressible, inviscid fluid in a uniform horizontal channel of finite depth. The variables in (1.1) and (1.2) are nondimensional and unscaled: *x* and *t* are proportional to position along the channel and time, respectively, and $\eta(x, t)$ and u(x, t) are proportional to the elevation of the free surface above the level of rest y = 0, and to the horizontal velocity of the fluid at a height $y = -1 + (1 + \eta(x, t))/\sqrt{3}$, respectively. (In terms of these variables the bottom of the channel is at y = -1.)

^{*} Corresponding author at: Department of Mathematics, University of Athens, 15784 Zographou, Greece. Tel.: +30 210 7276311; fax: +30 210 7276398.

E-mail address: doug@math.uoa.gr (V.A. Dougalis).

^{0378-4754/\$36.00} @ 2011 IMACS. Published by Elsevier B.V. All rights reserved. doi:10.1016/j.matcom.2011.09.006